1. The operators of position \hat{x} and of momentum \hat{p} act as follows on a wave function $\psi(x)$ in the position representation:

$$\hat{x}\psi(x) = x\psi(x), \,\hat{p}\psi(x) = -i\hbar\frac{d}{dx}\psi(x).$$

- *a*) How do \hat{x} and \hat{p} act on a wave function $\Phi(p)$ in the momentum representation?
- *b*) Show that the commutator $[\hat{x}, \hat{p}]$ is the same in position and momentum representation.
- *c)* Are these two operators

$$\hat{O}_1 = \hat{x}\hat{p}, \ \hat{O}_2 = \hat{x}^2\hat{p} - i\hbar\hat{x}$$

Hermitian or not? Motivate your answer.

2. Consider a Hamiltonian \hat{H} with eigenvalues $E_0 \leq E_1 \leq E_2 \cdots$. The variational principle which we discussed in class says that the ground state energy E_0 has the *upper bound* $U = \langle \phi | \hat{H} | \phi \rangle$, where $| \phi \rangle$ is an arbitrary "trial" wave function, normalized

by $\langle \phi | \phi \rangle = 1$.

• *a*) Expand the trial wave function $|\phi\rangle = \sum_{n=0}^{\infty} c_n |\psi_n\rangle$ in terms of the eigenfunctions $|\psi_n\rangle$ at energy E_n of \hat{H} and derive that $U = \sum_{n=0}^{\infty} |c_n|^2 E_n$. Why does this imply that $E_0 \leq U$?

We will now derive a different principle, which says that \hat{H} has at least one eigenvalue E_p in the interval

$$U - \sqrt{V - U^2} \le E_p \le U + \sqrt{V - U^2},\tag{1}$$

where $V = \langle \phi | H^2 | \phi \rangle = \sum_{n=0}^{\infty} |c_n|^2 E_n^2$.

• *b*) Prove that

$$V - U^2 = \sum_{n=0}^{\infty} |c_n|^2 (E_n - U)^2.$$

• *c)* Denote by E_p the eigenvalue of \hat{H} that is closest to *U*. Derive that

$$V - U^2 \ge (E_p - U)^2$$

and show that this implies the two inequalities in equation (1). Why is it wrong to conclude that the ground state energy has the *lower bound* $U - \sqrt{V - U^2}$?

continued on second page

3. The bosonic annihilation operator \hat{a} and creation operator \hat{a}^{\dagger} can be used to describe the state of photons at a given frequency. A laser emits photons in a socalled coherent state, given by

$$|\beta\rangle = e^{-|\beta|^2/2} e^{\beta \hat{a}^{\dagger}} |0\rangle.$$

Here β is a complex number and $|0\rangle$ is the vacuum state.

- *a*) Show that the coherent state is an eigenstate of the annihilation operator: $\hat{a}|\beta\rangle = \beta|\beta\rangle$.
- *b*) Show that two coherent states |α⟩ and |β⟩ are not orthogonal, by deriving that

$$|\langle \alpha | \beta \rangle|^2 = \exp(-|\alpha - \beta|^2).$$

- *c*) Calculate the first two moments of the photon number: $\bar{n} = \langle \beta | \hat{a}^{\dagger} \hat{a} | \beta \rangle$ and $\bar{n}^2 = \langle \beta | (\hat{a}^{\dagger} \hat{a})^2 | \beta \rangle$, and the variance var $n = \bar{n}^2 (\bar{n})^2$. Discuss how your result agrees with the expected Poisson statistics of photons in a coherent state.
- 4. We consider the Hamiltonian $\hat{H} = \hat{p}^2/2m + V(\hat{x})$ of a particle of mass m moving along the x-axis in a confining potential V(x). The eigenvalues of \hat{H} form the discrete spectrum E_0, E_1, E_2, \ldots We define the density of states $\rho(E) = \sum_{n=0}^{\infty} \delta(E E_n)$ and its Fourier transform

$$F(t) = \int_{-\infty}^{\infty} \rho(E) e^{-iEt/\hbar} dE = \sum_{n=0}^{\infty} e^{-iE_nt/\hbar}.$$

The dynamics from position x_0 to x_1 in a time *t* is described by the propagator

$$G(x_1, x_0; t) = \langle x_1 | e^{-i\hat{H}t/\hbar} | x_0 \rangle.$$

• *a*) Derive the following relation between F(t) and the integral of the propagator for equal initial and final position:

$$\int_{-\infty}^{\infty} G(x,x;t)dx = F(t).$$

Feynman showed that the propagator $G(x_1, x_0; t)$ can be written as an integral over all paths x(t') with $x(0) = x_0$ and $x(t) = x_1$,

$$G(x_1, x_0; t) = \sqrt{\frac{m}{2\pi i\hbar t}} \int_{x(0)=x_0}^{x(t)=x_1} \mathcal{D}[x(t')] e^{iS[x(t')]/\hbar}.$$

- *b*) What is the definition of *S*[*x*(*t'*)], how does the potential *V*(*x*) enter in this definition?
- *c)* Suppose that V(x) is a square well potential, so V(x) = 0 for 0 < x < W, while $V(x) \rightarrow \infty$ for x < 0 and x > W. Draw in an x-t diagram a path x(t) that contributes predominantly to the density of states $\rho(E)$ in the limit $\hbar \rightarrow 0$.