

Write NAME, INITIALS and STUDENT NUMBER on every sheet you hand in. Start each new problem on a new page.

All problems count for the same number of points (20) in the grading.

- 1) Consider a harmonic oscillator:  $[\hat{a}, \hat{a}^\dagger] = \hat{1}$ ,  $\hat{n} \equiv \hat{a}^\dagger \hat{a}$ ,  $\hat{n}$ -basis  $\{|n\rangle\}$  with  $n = 0, 1, 2, \dots$ ,  $\hat{a}|n\rangle = \sqrt{n}|n-1\rangle$ ,  $\hat{a}^\dagger|n\rangle = \sqrt{n+1}|n+1\rangle$ . An ensemble of such oscillators happens to be characterised by the following state (or: density) operator:

$$\hat{\rho} = \frac{1}{2}|2\rangle\langle 2| - \frac{i}{2}|2\rangle\langle 3| + \frac{i}{2}|3\rangle\langle 2| + \frac{1}{2}|3\rangle\langle 3|$$

- (a) Determine the state (or: density) matrix in the  $n$ -representation,  $\rho_{nn'}$ .
- (b) Does the given state operator describe a quantum system in a *pure state* or does it describe a *mixture*? Give an argument for your answer.
- (c) In case the state operator is  $\hat{\rho}$ , determine the following expectation values (or: ensemble averages):  $\langle a \rangle$ ,  $\langle a^\dagger \rangle$  and  $\langle n \rangle$ .
- 2) Consider a spin  $\frac{1}{2}$  object. Let the Hamilton operator be:  $\hat{H} = -\hbar\omega\hat{\sigma}_y$ . The normalised eigenstates of  $\hat{\sigma}_y$  can be used as a basis and are given by:

$$|+, y\rangle = \frac{1}{\sqrt{2}}(|+\rangle + i|-\rangle) \quad \text{and} \quad |-, y\rangle = \frac{1}{\sqrt{2}}(i|+\rangle + |-\rangle),$$

in terms of the eigenstates  $|+\rangle$  and  $|-\rangle$  of  $\hat{\sigma}_z$ .

- (a) Determine the matrix corresponding to  $\hat{\sigma}_z$  in the  $\sigma_y$ -representation (i.e. in the basis of eigenstates of  $\hat{\sigma}_y$ ).
- (b) Argue that the (unitary) matrix  $U$  that transforms a general state vector in the  $\sigma_z$ -representation into the corresponding state vector in the  $\sigma_y$ -representation is given by:

$$U = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -i \\ -i & 1 \end{pmatrix}.$$

Now, the state vector in the  $\sigma_z$ -representation at  $t = 0$  is:  $\chi(t = 0) = \begin{pmatrix} \cos \gamma \\ \sin \gamma \end{pmatrix}$ .

Call the state vector in de  $\sigma_y$ -representation:

$$\eta(t) = \begin{pmatrix} e(t) \\ f(t) \end{pmatrix}.$$

- (c) Determine  $\eta(t = 0)$  and solve the Schrödinger-equation  $i\hbar\frac{\partial}{\partial t}\eta(t) = \hat{H}\eta(t)$  in the  $\sigma_y$ -representation.
- (d) Using the results derived in (a) and (c), show that the expectation value of  $\hat{\sigma}_z$  at time  $t$  is given by:

$$\langle \sigma_z \rangle(t) = \cos(2\omega t - 2\gamma).$$

- 3) Consider a single species of bosons with annihilation- and creation operators  $\hat{a}$  and  $\hat{a}^\dagger$ , respectively. The Hamilton operator for this quantum many-body system is:

$$\hat{H} = \omega \left( \hat{a}^\dagger \hat{a} + \frac{1}{2} \right) + \frac{1}{2} \Delta \left( \hat{a}^\dagger \hat{a}^\dagger + \hat{a} \hat{a} \right) . \quad (1)$$

We take  $\hbar = 1$  throughout this problem. The following transformation is useful to gain insight into the properties of this quantum system:

$$\hat{b} = \lambda \hat{a} + \mu \hat{a}^\dagger, \quad (2a)$$

$$\hat{b}^\dagger = \lambda^* \hat{a}^\dagger + \mu^* \hat{a}, \quad (2b)$$

where  $\lambda$  and  $\mu$  are complex numbers.

- (a) Show that the transformation (2) preserves the canonical commutation relations provided  $|\lambda|^2 - |\mu|^2 = 1$ .
- (b) Assuming  $\lambda$  and  $\mu$  to be real and using the result of (a), show that transformation (2) brings the Hamiltonian (1) into the form:

$$\hat{H} = \tilde{\omega} \left( \hat{b}^\dagger \hat{b} + \frac{1}{2} \right) . \quad (3)$$

Provide expressions for  $\tilde{\omega}$ ,  $\lambda^2$ , and  $\mu^2$  in terms of  $\omega$  and  $\Delta$ .

- (c) If the bosons characterized by  $\hat{a}$  and  $\hat{a}^\dagger$  are considered as excitations of a harmonic oscillator, the connection with the (Hermitian) operators for position and momentum of the oscillator,  $\hat{x}$  and  $\hat{p}$ , is given by:

$$\hat{a} = \frac{1}{\sqrt{2}} \left( \sqrt{m\omega} \hat{x} + \frac{i \hat{p}}{\sqrt{m\omega}} \right) . \quad (4)$$

Express the Hamiltonian in terms of  $\hat{x}$  and  $\hat{p}$  for the special case  $\Delta = \omega$ . How would you interpret this result physically?

- 4) Consider a many-body system corresponding to  $N$ -particle systems consisting of  $N$  identical spin- $\frac{1}{2}$  particles with Hamilton operator:

$$\hat{H}_N = \sum_{i=1}^N \hat{h}^{(i)} \quad , \quad \hat{h} = \frac{\hat{p}^2}{2m} + B \hat{\sigma}_x .$$

Use the discrete  $\vec{k}\sigma$ -representation ( $\hat{p}|\vec{k}\sigma\rangle = \hbar\vec{k}|\vec{k}\sigma\rangle$ ,  $\hat{\sigma}_z|\vec{k}\sigma\rangle = \sigma|\vec{k}\sigma\rangle$ ,  $\sigma = +1, -1$ ).

- (a) Give the fundamental algebraic relations of the annihilation- and creation-operators  $\hat{a}_{\vec{k}\sigma}$  and  $\hat{a}_{\vec{k}\sigma}^\dagger$ .
- (b) Show that the many-body energy operator  $\hat{H}$  in the  $\vec{k}\sigma$ -representation is of the following form and determine  $f$ ,  $g_+$  and  $g_-$ :

$$\hat{H} = \sum_{\vec{k}\sigma} f \hat{a}_{\vec{k}\sigma}^\dagger \hat{a}_{\vec{k}\sigma} + \sum_{\vec{k}} \left( g_+ \hat{a}_{\vec{k},1}^\dagger \hat{a}_{\vec{k},-1} + g_- \hat{a}_{\vec{k},-1}^\dagger \hat{a}_{\vec{k},1} \right) .$$

- (c) Compute  $[\hat{H}, \hat{a}_{\vec{k}\sigma}]$ .
- (d) Calculate the Heisenberg-picture operator  $\hat{c}_{\vec{k}}(t)$  with  $\hat{c}_{\vec{k}} \equiv \hat{a}_{\vec{k},1} + \hat{a}_{\vec{k},-1}$ .
- 5) In classical mechanics a harmonic oscillator with mass  $m$  and (angular) frequency  $\omega$  in one spatial dimension is given by the following Lagrangian  $L$  (the dot above  $x$  means differentiation with respect to time):

$$L(\dot{x}, x) = \frac{1}{2}m\dot{x}^2 - \frac{1}{2}m\omega^2 x^2 \quad .$$

The equation of motion follows from the Euler-Lagrange equation:

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{x}} - \frac{\partial L}{\partial x} = 0 \quad .$$

The action  $\mathcal{S}$  for the path  $x(t)$  is given by:

$$\mathcal{S}[x(t)] = \int dt L(\dot{x}, x) \quad ,$$

where the integral runs from begin- to end-time of the path.

- (a) Calculate the action  $\mathcal{S}_{cl}$  for the classical path of the harmonic oscillator that starts at time  $t = 0$  at position  $y$  and ends at time  $t'$  with *velocity* equal to:  $-y\omega \sin(\omega t')$ .
- (b) The quantummechanical *propagator* for a harmonic oscillator that at time  $t_0$  has position  $x_0$  and at time  $t'$  has position  $x'$  is given by:

$$\langle x', t' | x_0, t_0 \rangle = \langle x' | e^{-i\hat{H}(t'-t_0)/\hbar} | x_0 \rangle \quad ,$$

where the Hamiltonian is given by:

$$\hat{H} = \frac{\hat{p}_x^2}{2m} + \frac{1}{2}m\omega^2 \hat{x}^2 \quad .$$

Show that the propagator for the harmonic oscillator with position  $x$  at  $t = 0$  and position  $x'$  a *short* time interval  $\Delta t$  later is of the following form and determine the function  $T(p, x, x')$ :

$$\langle x' | e^{-i\hat{H}\Delta t/\hbar} | x \rangle = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} dp \exp \{ iT(p, x, x')\Delta t/\hbar \} + \mathcal{O}(\Delta t)^2 \quad .$$

If necessary, use that the scalar product of eigenstates of  $\hat{x}$  en  $\hat{p}_x$  is given by:

$$\langle x | p \rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{ipx/\hbar} \quad .$$

