



# Quantum Theory

## lecture 2: Symmetry

### the force of symmetry

- conservation laws S-4.1
- unitary & anti-unitary symmetries S-4.4
- parity S-4.2
- time-reversal & Kramers degeneracy S-4.4
- Galilean invariance

S = Sakurai (2nd edition)

# conservation laws

Hamiltonian invariant under a unitary transformation

$$\hat{U}\hat{H}\hat{U}^\dagger = \hat{H} \Rightarrow [\hat{U}, \hat{H}] = 0$$

$$\hat{U} = e^{i\hat{A}} \Rightarrow i\hbar \frac{d\hat{A}}{dt} = [\hat{A}, \hat{H}] = 0$$

observable corresponding to  $\hat{A}$  is conserved

(the Hermitian operator  $\hat{A}$  is called the generator of the unitary symmetry  $\hat{U}$ )

**translational symmetry:**  $\hat{U} = e^{i\hat{p}a/\hbar} \Rightarrow$  momentum conserved

see exercise 1.2, or consider an infinitesimal translation over  $a$ :

$$\delta\hat{U} = 1 + \delta a \partial/\partial x \text{ with } \delta a = a/N$$

then compose  $N$  such translations and take the limit  $N \rightarrow \infty$

$$\hat{U} = \left(1 + \frac{a}{N} \frac{\partial}{\partial x}\right)^N = e^{a\partial/\partial x} = e^{ia\hat{p}/\hbar}$$

# unitary & anti-unitary symmetries

Wigner's theorem: every symmetry  $S$  is either a **unitary** operator or an **anti-unitary** operator.

Symmetry:  $|\langle S\psi | S\chi \rangle|^2 = |\langle \psi | \chi \rangle|^2$  for all states  $\psi, \chi$ .

Unitary case:  $S = U, UU^\dagger = U^\dagger U = 1,$

Anti-unitary case:  $S = UK,$  with  $K =$  complex conjugation

for example, time-reversal:  $\psi^*(\mathbf{x}, t) = \psi(\mathbf{x}, -t)$

Sketch of a proof for a 2D Hilbert space (spin-1/2):

$|\psi\rangle$  is associated with a vector  $\vec{n}$  on the Bloch sphere

$$|\langle \psi | \chi \rangle|^2 = \frac{1}{2}(1 + \vec{n}_1 \cdot \vec{n}_2).$$

symmetry is angle-preserving mapping of the unit sphere on itself:  
only rotations ( $S = U$ ) or rotations + reflection ( $S = UK$ ).

# parity (= spatial inversion)

see exercise 2.1

$$\mathcal{P}\hat{q}\mathcal{P}^{-1} = -\hat{q} \quad \text{and} \quad \mathcal{P}\hat{p}\mathcal{P}^{-1} = -\hat{p}, \quad \mathcal{P}^2 = 1$$

unitary operator, eigenvalues  $\pm 1$  (odd or even parity)

# time-reversal

see exercise 2.2

$$\mathcal{T}\hat{q}\mathcal{T}^{-1} = \hat{q} \quad \text{and} \quad \mathcal{T}\hat{p}\mathcal{T}^{-1} = -\hat{p}, \quad \mathcal{T}^2 = \pm 1.$$

anti-unitary operator, examples:  $\mathcal{T}\Psi = \Psi^*$  or  $\mathcal{T}\Psi = \sigma_y\Psi^*$

# Kramers degeneracy

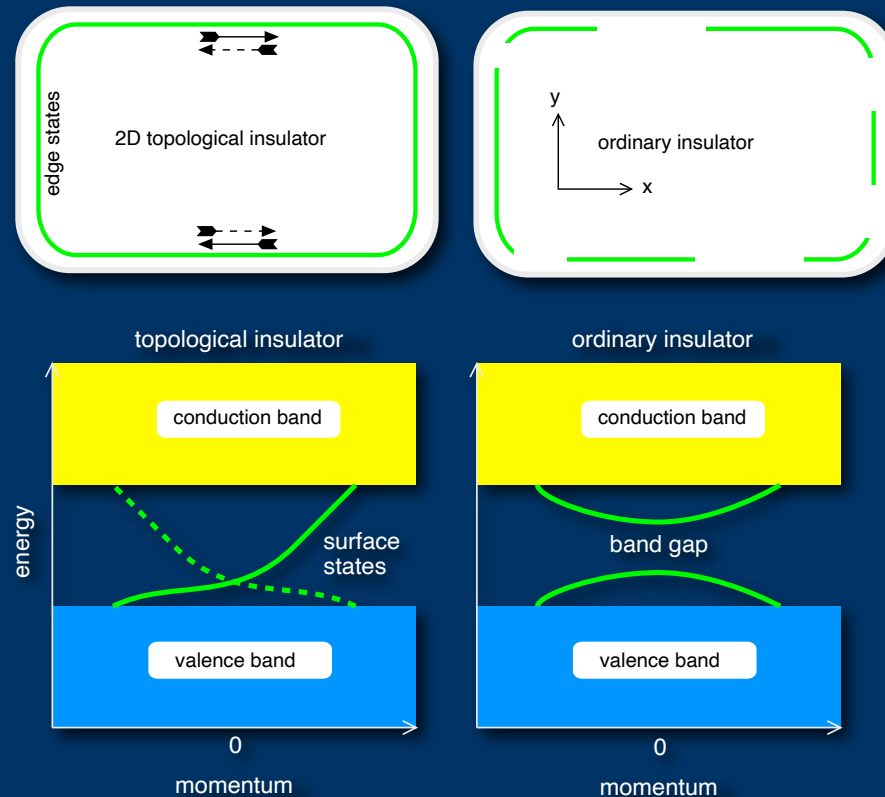
if  $\mathcal{H}$  commutes with  $\mathcal{T}$  and  $\mathcal{T}^2 = -1$ , then for every eigenvalue there are at least two independent eigenvectors

if  $\mathcal{H}|\psi\rangle = E|\psi\rangle$  then also  $\mathcal{H}|\psi'\rangle = E|\psi'\rangle$  with  $|\psi'\rangle = \mathcal{T}|\psi\rangle$ ;

use  $\mathcal{T}^2 = -1$  to prove by contradiction that there cannot be a  $\lambda \neq 0$  such that  $|\psi\rangle = \lambda|\psi'\rangle$

# quantum spin Hall effect

Charles Kane & Gene Mele (2004) — first observation by Laurens Molenkamp (2007)



Kramers degeneracy protects the crossing at zero momentum;  
mixing of left-movers and right-movers is forbidden

Hellmann-Feynman:  $\langle v \rangle = \langle \psi | \partial H / \partial p | \psi \rangle = dE / dp$